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Direct current stimulation boosts hebbian plasticity in vitro

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ABSTRACT

Background: There is evidence that transcranial direct current stimulation (tDCS) can improve learning performance. Arguably, this effect is related to long term potentiation (LTP), but the precise biophysical mechanisms remain unknown.

Hypothesis: We propose that direct current stimulation (DCS) causes small changes in postsynaptic membrane potential during ongoing endogenous synaptic activity. The altered voltage dynamics in the postsynaptic neuron then modify synaptic strength via the machinery of endogenous voltage-dependent Hebbian plasticity. This hypothesis predicts that DCS should exhibit Hebbian properties, namely pathway specificity and associativity.

Methods: We studied the effects of DCS applied during the induction of LTP in the CA1 region of rat hippocampal slices and using a biophysical computational model.

Results: DCS enhanced LTP, but only at synapses that were undergoing plasticity, confirming that DCS respects Hebbian pathway specificity. When different synaptic pathways cooperated to produce LTP, DCS enhanced this cooperation, boosting Hebbian associativity. Further slice experiments and computer simulations support a model where polarization of postsynaptic pyramidal neurons drives these plasticity effects through endogenous Hebbian mechanisms. The model is able to reconcile several experimental results by capturing the complex interaction between the induced electric field, neuron morphology, and endogenous neural activity.

Conclusions: These results suggest that tDCS can enhance associative learning. We propose that clinical tDCS should be applied during tasks that induce Hebbian plasticity to harness this phenomenon, and that the effects should be task specific through their interaction with endogenous plasticity mechanisms. Models that incorporate brain state and plasticity mechanisms may help to improve prediction of tDCS outcomes.

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Introduction

Transcranial direct current stimulation (tDCS) studies in humans have recently exploded in number and scope [1—4]. While these studies have seen varying degrees of success [1], in aggregate they suggest that stimulation with weak constant current can have long term effects on cognitive function [5]. One of the predominant theories to explain these long term effects is that stimulation affects synaptic plasticity [6], although a variety of alternatives have also been proposed [7,8] and are being explored [9]. The synaptic

Abbreviations: tDCS, transcranial direct current stimulation; DCS, direct current stimulation; LTP, long term potentiation; TBS, theta burst stimulation; STDP, spike-timing dependent plasticity.

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plasticity theory is consistent with an array of findings from pharmacological studies in humans [10] as well as animal electrophysiology studies conducted in-vivo [9,11,12] and in-vitro [13–16]. However, the biophysical mechanism for such plasticity effects is unknown.

Polarization of neuronal membranes in response to extracellular electric fields has been well characterized [17–23], as has the membrane potential-dependence of Hebbian plasticity [24–27]. While it is straightforward to draw a connection between these phenomena, their interaction can be complex. For example, we previously observed that the effects of DCS depend on both the location of active synapses and the precise temporal patterns of activity used to induce plasticity [14]. These results suggest that the effects of DCS depend on the interaction between the induced electric field, neuron morphology, and the endogenous brain dynamics. Given this complexity, it is perhaps no surprise that results

from human clinical trials with tDCS have remained inconclusive [28–31], or that optimization of tDCS protocols has been slow. For example, there is an ongoing debate as to whether tDCS should be applied before, during, or after a behavioral or cognitive task [32–34].

We propose that DCS causes small changes in postsynaptic membrane potential during ongoing endogenous synaptic activity. The altered voltage dynamics in the postsynaptic neuron then modify synaptic strength via the machinery of endogenous voltage-dependent Hebbian plasticity. An implication of this hypothesis is that the effects of DCS should exhibit similar properties as the endogenous Hebbian plasticity that it is paired with. Two of these properties, pathway specificity and pathway associativity [19,20], support functionally specific learning of cell assemblies in neural networks [21,22]. tDCS may therefore enhance functionally specific learning by acting through this Hebbian mechanism.

We induced LTP in hippocampal brain slices using theta rhythms (theta burst stimulation, TBS), and confirm that this form of "endogenous" plasticity is pathway specific and associative. Applying DCS during plasticity induction boosted the amount of LTP, while maintaining the pathway-specific and associative properties of the underlying endogenous plasticity. Additional experiments and computer simulations support the hypothesized model in which DCS achieves these effects through altered neuronal excitability and subthreshold depolarization in dendrites during ongoing synaptic input.

We present what is, to our knowledge, the first computational model of the effects of DCS on synaptic plasticity, which reconciles several experimental results. The model makes specific and testable predictions for both how tDCS should alter plasticity when paired with various endogenous brain states, and how this can inform the design of tDCS protocols. Specifically, the most effective tDCS interventions should be those that pair stimulation concurrently with behavioral training and that performance gains should be specific to the learned task.

Results

Anodal DCS boosts LTP

To mimic learning during a training task we induced LTP by applying TBS in the hippocampal Schaffer collateral pathway (4 pulses at 100 Hz repeated for 15 bursts at 5 Hz, 3 s total). We applied acute anodal or cathodal DCS (see Methods) for the duration of the LTP-induction protocol (20 V/m; Fig. 1A). When paired with anodal DCS, the resulting LTP was increased compared to TBS alone (Fig. 1B; control: 1.287+-0.025, N = 52 slices; anodal: 1.397+-0.047, N = 32 slices, p = 0.027). However, cathodal stimulation had no significant effect (Fig. 1B; cathodal: 1.243+-0.031, N = 12 slices, p = 0.424).

Electric field interacts with plasticity induction on millisecond timescale

Membrane polarization during DCS has been well characterized [17–22] and is well described by cable models of stimulated neurons (Fig. 1A) [21–23]. We previously argued that the effects of DCS on tetanus-induced LTP are due to membrane polarization [14]. If this is the case for TBS-induced LTP as well, then there is no need for the DCS to be constant over long periods of time. It would suffice for the DCS field to coincide with TBS synaptic inputs on the time scale of the neuronal membrane time constant (e.g. 30 ms) [18]. To test for this, we applied theta-frequency alternating current stimulation (sinusoidal 5 Hz at 20 V/m) during TBS induction. The peak phase of this alternating current corresponds to the same electric field as

anodal DCS, while trough corresponds to cathodal DCS. When TBS bursts were timed to coincide with the peak of the alternating current, LTP was enhanced, as with anodal DCS (Fig. 1C; control: 1.287+-0.025, N=52; peak: 1.467+-0.093, N=9, p=0.014; N here and below indicates the number of slices). TBS timed to the trough of the alternating current had no significant effect on LTP, as with cathodal DCS (Fig. 1C; trough: 1.184+-0.035, N=6, p=0.173). These data suggest that the electric field need only coincide with potentiating synaptic input on the millisecond timescale, and does not require any prolonged buildup of DCS effects in order to affect LTP. This is consistent with the notion that instantaneous membrane polarization due to DCS is what interacts with synaptic activity to modulate the resulting plasticity [14].

Effect of DCS on LTP is pathway specific

Hebbian synaptic plasticity is classically characterized as a pathway specific process, i.e. only pathways that are coactive with the postsynaptic neuron are strengthened [35]. Our proposal that DCS enhances LTP through membrane potential implies that the effects of DCS should follow this pathway specificity. We tested this by monitoring two independent synaptic pathways in CA1 (Fig. 2A). During induction, the strong pathway received TBS while the other inactive pathway was not stimulated. As expected, LTP was observed in the strong pathway (Fig. 2B black; 1.377+-0.052, N = 16, p = 2.8E-6), but not the inactive pathway (Fig. 2B gray; 0.986+-0.031, N = 14, p = 0.657), demonstrating the wellestablished pathway specificity of LTP [35]. When this induction protocol was paired with anodal DCS, LTP was enhanced only in the strong pathway (Fig. 2B red; 1.613+-0.071, N = 14, p = 0.011 vs. control), while the inactive pathway was unaffected (Fig. 2B light red; 0.971+-0.028, N = 14, p = 0.724 vs. control), showing that this effect of DCS is specific to the potentiated pathway.

DCS boosts hebbian associativity

Another important property of Hebbian plasticity is pathway associativity, which is a cellular mechanism thought to underlie the formation of cell assemblies and associative learning [35–37]. Pathway associativity refers to the potentiation of separate synaptic pathways arriving onto the same postsynaptic neuron when they cooperate to drive the postsynaptic cell. For example, a synaptic input that is too weak on its own to induce plasticity can undergo plasticity if it is coactivated with a strong input that helps to drive the postsynaptic cell.

We tested how DCS affects Hebbian associativity by again monitoring two synaptic pathways. First, only a weak input (15 pulses at 5 Hz) was used during induction (Fig. 3A). In the absence of DCS, no lasting plasticity was observed in this weakly activated pathway (Fig. 3A gray; 0.998+-0.041, N = 13, p = 0.966) or the other inactive pathway (Fig. 3A black; 0.958+-0.037, N = 13, p = 0.275). DCS also had no effect on the weak (Fig. 3A light red; 1.041+-0.038, N = 13, p = 0.445) or inactive pathway (Fig. 3A red; 0.963+-0.011, N = 13, p = 0.908). This result further confirms the specificity of DCS effects, in that pathways that are not undergoing plasticity are unaffected by DCS.

In a second experiment, the weak input is now paired with a strong input (TBS) during induction (Fig. 3B). During induction, weak pathway inputs are timed to arrive at precisely the same time as the second pulse of each theta burst. This induces LTP in the strong pathway as before (Fig. 3B black; 1.435+-0.067, N = 13, p = 3.1E-5), but now the weak pathway is also potentiated (Fig. 3B gray; 1.115+-0.031, N = 13, p = 0.003), replicating classic associativity between the two pathways [35]. If this protocol is paired with DCS during induction, LTP is now boosted in both the strong (Fig. 3B

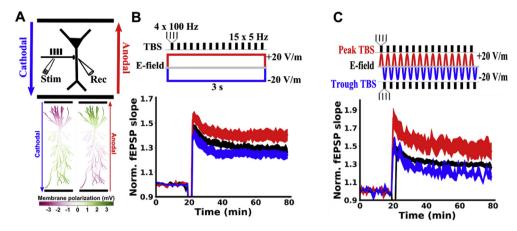


Fig. 1. Soma-depolarizing electric fields enhance TBS-induced LTP in hippocampal Schaffer Collateral pathway. A) Top: Schematic of the experimental setup, showing the orientation of anodal (red) and cathodal (blue) electric fields generated by parallel wires (black horizontal lines). Location of stimulation (Stim) with TBS and recording (Rec) of field excitatory postsynaptic potentials (fEPSP) are indicated relative to a CA1 pyramidal neuron soma (black triangle). Bottom: Membrane polarization throughout a model pyramidal neuron in response to 20 V/m anodal (red) or cathodal (blue) DCS. Green compartments are depolarized due to DCS, while magenta compartments are hyperpolarized by DCS. B) Constant current stimulation applied during TBS modulates the resulting LTP measured as a change in fEPSP normalized to baseline. C) Alternating current stimulation (5 Hz) was applied and TBS bursts were timed to either the peak (red) or the trough (blue) of the sinusoidal alternating current. Note that the applied electric field at the peak of the alternating current is identical to anodal constant current, as is the case for the trough of the alternating current and cathodal constant current. The effects of alternating currents are similar to those of the analogous constant current paradigm, indicating that plasticity modulation is consistent with the instantaneous incremental membrane polarization on a millisecond timescale. LTP induction is applied at the 20 min mark. All data are normalized to the mean of the 20 baseline responses before induction and are represented as mean ± s.e.m across slices. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

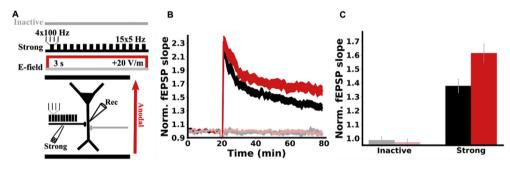


Fig. 2. DCS effect is specific to the potentiated pathway. A) Schematic of the experimental setup. Two synaptic pathways are monitored before and after plasticity induction. During induction, one pathway is activated with TBS (black, strong), while the other pathway is inactive (gray), and anodal DCS is applied across the slice throughout the duration of induction (3 s, red). **B)** Plasticity is pathway specific and so are DCS effects. LTP was observed only in the pathway that received TBS (black trace), demonstrating pathway specificity. Anodal DCS enhanced LTP only in the potentiated pathway (red vs black) and had no effect on the inactive pathway (light red vs. gray), upholding Hebbian specificity. fEPSP slopes normalized to the mean of the 20 of baseline responses prior to induction is applied at the 20 min mark. **C)** Summary of pathway specific effects of DCS. The mean of the last 10 normalized slopes (51–60 min after induction) are used for each slice. Data are represented as mean ± s.e.m across slices. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

red c. f. black; 1.705+-0.094, N = 13, p = 0.029) and the weak pathway (Fig. 3B light red c. f. gray; 1.242+-0.029, N = 13, p = 0.006). DCS therefore enhances the Hebbian associativity between the strong and weak pathways (Fig. 3D). We note that plasticity was similar in the strong (TBS) pathway, regardless of whether it was paired with the weak pathway (Fig. 3C black), and that the effect of DCS on the strong pathway was indifferent to pairing as well (Fig. 3C red).

Effects are consistent with DCS modulation of somatic spiking

We hypothesized that the effects of DCS on TBS-induced LTP are due to membrane polarization. However, DCS will alter membrane potential in both the soma and apical dendrites of pyramidal neurons, but with opposite polarities [18,23]. We therefore aimed to test whether the effects of DCS on LTP were consistent with somatic or dendritic membrane polarization. To do so, we took a similar approach as in previous work [14]. LTP was induced by stimulation of Schaffer collaterals with TBS in either apical or basal dendritic compartments of CA1 (Fig. 4B). DCS is expected to have opposite

effects on dendritic membrane potential in basal as opposed to apical dendrites (Fig. 5A) [18,23]. Effects due to DCS-induced *dendritic* polarization should therefore be *opposite* when synapses are activated in apical or basal dendrites. However, effects due to DCS-induced *somatic* polarization should be the *same*, regardless of the location of synaptic activation (i.e. there is only one soma per neuron). Therefore, observing different effects in apical and basal compartments would rule out somatic polarization as a main determinant of the plasticity modulation.

Here we found that DCS had the same effect on LTP in both basal and apical dendrites (Fig. 4C and D). This result is consistent with plasticity effects of DCS being driven primarily by effects on somatic spiking. To further test this, we looked at measures of dendritic integration and somatic spiking in each condition (Fig. 4A, see Methods for details of analysis). Indeed, we found that DCS had a similar effect on somatic spiking (Fig. 4E), but opposite effects on dendritic integration in apical versus basal dendrites (Fig. 4F). Thus, the effect of DCS polarity on LTP mirrors that of the effect on the soma, but not dendrites.

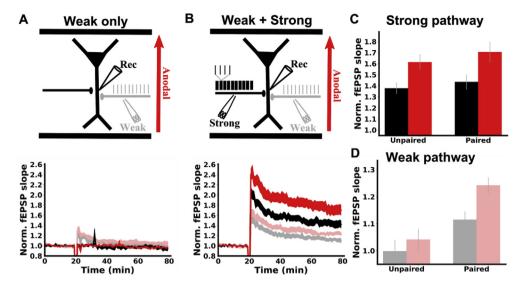


Fig. 3. DCS enhances associativity between synaptic pathways. A) Top: schematic of experimental design. Two synaptic pathways were monitored. During induction, one pathway was weakly activated at 5 Hz with 15 pulses (gray), while the other pathway was inactive (black). Anodal DCS was applied throughout induction (3 s, red). Bottom: weak synaptic activation had no lasting effect on synaptic synaptic strength in either pathway with DCS (red, light red) or without DCS (gray, black). B) Top: schematic of experimental design. Again, two synaptic pathways were monitored. Now during induction, one pathway was activated with a TBS protocol (strong, black). The other pathway was activated with 15 pulses at 5 Hz (weak, gray). Weak pathway pulses were temporally aligned to the second pulse in each TBS burst. Bottom: without DCS, the strong pathway was potentiated (black) and the weak pathway was now also potentiated (gray), demonstrating associative plasticity between these pathways. With DCS, LTP was enhanced in the strong pathway (red) and the weak pathway (light red), demonstrating that the associativity between pathways was enhanced. C) Summary of LTP experiments in the strong pathway. Pairing with the weak pathway did not increase strong pathway. LTP, and DCS had a similar effect on LTP in both cases. D) Summary of LTP experiments in the weak pathway. EPSP slopes are normalized to the mean of the 20 of baseline responses prior to induction. Induction is applied at the 20 min mark in panels A,B. The mean of the last 10 normalized slopes (51–60 min after induction) are used for each slice in panels C,D. Data are represented as mean ± s.e.m. across slices. . (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

Computational model

To further understand how changes in membrane potential due to DCS lead to the observed changes in plasticity, we turned to a computational model. We modeled a CA1 pyramidal neuron based on a previously validated biophysical model, using the NEURON software package [38-40]. To simulate the effects of DCS, we applied a uniform extracellular electric field (voltage gradient) with NEURON's extracellular mechanism [23]. This extracellular field is known to polarize the cellular membrane with opposite polarities in apical and basal compartment (Fig. 5A) [18]. To calculate activitydependent synaptic plasticity, we used a voltage-based plasticity rule [41] that has been used previously to replicate a wealth of synaptic plasticity data [41–43]. Here we manually selected parameters for this plasticity rule such that we could qualitatively reproduce canonical spike-timing dependent plasticity (STDP) experiments [26,43](Supplemental Fig. S1) and the effects of DCS on synaptic plasticity in our own TBS experiments (compare experiments of Fig. 4D-F with model results of Fig. 5). The model also reproduces the experimental results with alternating current stimulation (compare experiment of Fig. 1C with model results of Fig. S3). All simulation results that follow use the same parameters unless specified otherwise (Supplemental Tables).

Associativity is enhanced through somatic spiking in simulations

Using the computational model, we then aimed to understand how DCS modulates TBS-induced LTP, while preserving specificity and associativity. Pathway specificity is explicitly built into the voltage-based plasticity rule of the model [41], following well established experimental results [35], namely synaptic weights are only allowed to change at active presynaptic inputs (see Methods). Since DCS does not by itself cause presynaptic activity, it cannot affect synaptic efficacy of the inactive pathway. Thus, the

incremental membrane polarization due to DCS upholds Hebbian synapse specificity.

It is less clear however, exactly how DCS is able to boost associativity between the weak and strong pathways (Fig. 3). We hypothesized that DCS boosted associativity through a boost of somatic spikes, which propagate to both weak and strong pathway synapses. To test this in the model, we simulated the experiments of Fig. 3, by activating one pathway with TBS (strong) and the other pathway with the 5 Hz stimulation (weak). When the weak pathway was activated alone no spikes were generated and only very weak plasticity was observed (Fig. 6D, weak only, black). Applying DCS in this case led to only minor changes in plasticity, as in our experiments (Fig. 6D, weak only, compare red and black). However, when the weak input was paired with the strong input, action potentials were generated in the soma that back-propagated to weak pathway synapses (Fig. 6B, black), and LTP was observed (Fig. 6D, weak + strong, black). Therefore, the weak and strong pathway become associated by cooperating to produce somatic spikes, which are then shared by both pathways.

When strong and weak pathways were paired, DCS facilitated the initiation of somatic spikes (Fig. 6B) and advanced their timing relative to the presynaptic input (Fig. 6C), due to increased depolarization of the soma. This led to a boost in weak pathway plasticity only when paired with the strong input (Fig. 6D, weak + strong), as observed experimentally (Fig. 6E; same as Fig. 3D).

To further validate the role of somatically initiated spikes in generating this DCS effect, we repeated the previous simulations, but set the voltage-gated sodium conductance to zero in the soma and axon (Fig. 6A bottom). This is analogous to the local application of TTX at the soma [44], preventing the initiation of spikes there. If the strength of synaptic stimulation is increased, spikes can still be generated, but they initiate locally in the dendrite (Fig. 6B bottom). Anodal DCS now reduces the probability of these spikes (Fig. 6B

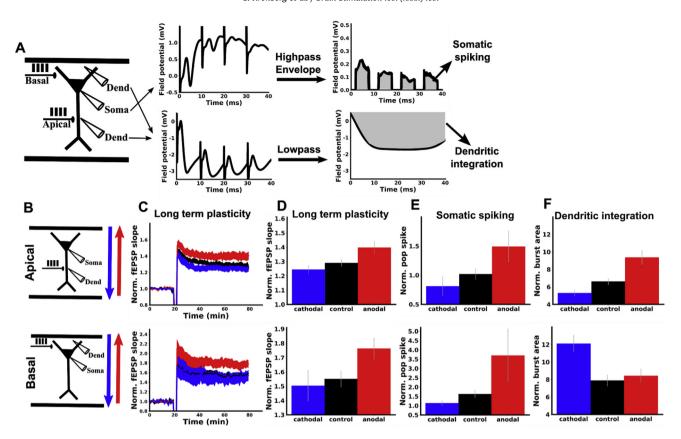


Fig. 4. DCS modulation of TBS-LTP is consistent with modulation of somatic spiking rather than dendritic integration. LTP was induced with TBS in either apical (top row, B—F) or basal (bottom row, B—F) dendritic regions of CA1. TBS induction was paired with anodal (red), cathodal (blue), or no DCS (black). A) Schematic of experiments and methods for deriving somatic and dendritic activity metrics. For both apical and basal protocols, one recording electrode was placed in the dendrites (Dend) near the bipolar stimulating electrode (Apical or Basal) and one electrode was placed near the CA1 somatic layer (Soma). Examples of raw voltage traces from each recording electrode during a single burst of the induction protocol are displayed in the middle panel. To derive a measure of dendritic integration, the dendritic recording was low-pass filtered, and the integral of this filtered signal was taken for each burst during TBS (gray area). To derive a measure of somatic population spiking, the somatic recording was high-pass filtered, and the integral of this signal's envelope during each burst was used (gray area; excludes periods of stimulation artefacts; see methods). B) Schematic of apical (top row) and basal (bottom row) experiments. C) Anodal DCS (red) boosts LTP in both and apical and basal dendrites compared to control (black). Cathodal DCS (blue) had no significant effect in either apical of basal dendrites. TBS was applied with or without DCS at the 20 min mark. Note that the top panel is identical to Fig. 1A (shown again here for comparison). D) Summary of the data in C. The mean of the last ten normalized responses were used for each slice. E) Population spiking measured for the first bipolar input pulse of the last burst (see Supplemental Fig. S2C for all bursts during induction). F) Population dendritic integration for the last burst of TBS (see Supplemental Fig. S2F for all bursts during induction). All data normalized to the reader is referred to the Web version of this article.)

bottom) and delayed their timing relative to the weak pathway input (Fig. 6C bottom), due to DCS-induced hyperpolarization of the apical dendrites. A prediction of this model is therefore that TTX applied locally at the soma, would cause anodal DCS to have the opposite effect on pathway associativity (Fig. 6D bottom), namely anodal DCS weakens rather than boosts LTP.

Taken together, the results of Fig. 6 suggest that DCS can enhance associativity by facilitating the initiation of somatic spikes. The additional spikes can spread to synapses in both pathways and increase LTP, leading to a stronger association between the pathways.

Interaction between synapse location and induction protocol

In a previous study, we used 20 Hz tetanic stimulation to induce LTP. We observed that a boost in LTP required opposite DCS polarities for apical and basal dendrites, suggesting that dendritic rather than somatic effects were dominant for this protocol [14], (Fig. 7B, top two rows). This appears inconsistent with the previous claim that DCS effects are mediated primarily through somatic spiking (Figs. 4—6). However, the computational model can readily reconcile these results if we consider the different endogenous

membrane voltage dynamics during 20 Hz tetanus and TBS protocols. For 20 Hz tetanic stimulation, inputs arriving early in the tetanus may elicit somatic spiking (Fig. 7D), but these inputs quickly become subthreshold due to short term synaptic depression (Fig. 7E). Since the majority of input pulses remain subthreshold, plasticity at these synapses is dominated by the local subthreshold dendritic potential (Fig. 7F). Because the DCS-induced polarization is opposite in apical and basal dendritic compartments, the effects on plasticity are also opposite there (Fig. 7B and C; compare 20 Hz apical to 20 Hz basal).

During TBS on the other hand, each burst in the induction is close to threshold at the soma. Somatic action potentials are generated throughout the induction, and plasticity at each synapse is dominated by the back-propagation of these spikes (Fig. 7D and E; bottom two rows). Effects of DCS on plasticity therefore follow the effects on somatic spike generation, regardless of the dendritic location of the synapses (Fig. 7B and C; compare TBS apical to TBS basal). Indeed, our experimental recordings of somatic spikes and dendritic integration in the CA1 population support this notion (Fig. 4). Performing a similar analysis in the model recapitulates this result (Fig. 5, c. f. Fig. 4D–F).

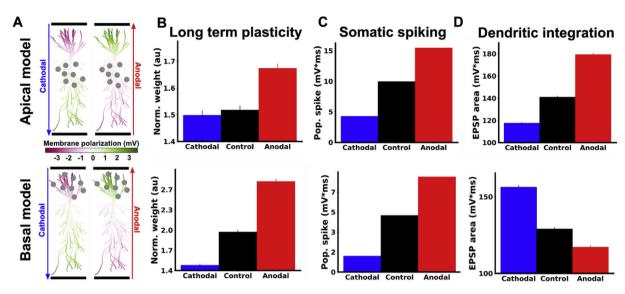


Fig. 5. Model captures the effects of DCS on long term potentiation, somatic spiking and dendritic integration. A) Membrane polarization throughout model pyramidal neuron in response to 20 V/m anodal (red) or cathodal (blue) DCS. Green compartments are depolarized due to DCS, while magenta compartments are hyperpolarized by DCS. Gray circles indicate the location of synapses in apical (top row) or basal (bottom row) compartments that are activated with TBS. B) Model predictions of changes in synaptic weights qualitatively match LTP experiments (c.f. Fig. 4D). The vertical axis (Norm. Weight) is the average weight of all activated synapses at the end of simulation, calculated offline using the learning rule [41]. C) Effects of DCS on somatic activity qualitatively match experimental measurements (c.f. Fig. 4E). The vertical axis is the average across all neuron somas of the high-pass filtered voltage envelope (see methods). D) Effects of DCS on dendritic integration qualitatively match experimental measurements (c.f. Fig. 4F). The vertical axis is the average across all recorded dendritic locations of the high-pass filtered envelope of the voltage (see methods). (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

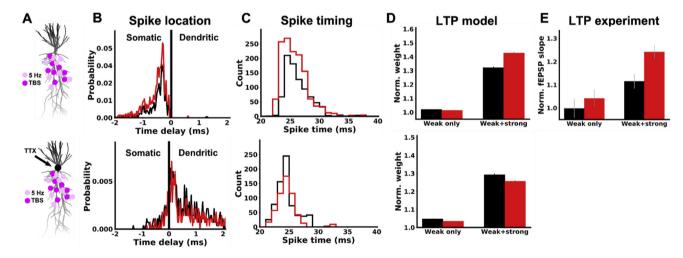


Fig. 6. Boost of associative LTP is also explained by the effect of DCS on somatic spikes in computational model. Top row: A) Simulated neuron morphology, showing an example of how synapses are distributed in the weak (5 Hz, light pink) and strong (TBS, magenta) pathways. B) Distribution of time delays between spikes observed in the soma and at weak pathway synapses for 20 V/m anodal stimulation (red) or control (black). Negative time delays correspond to spikes that occur in the soma first. Due to variable propagation delays between synapses, it is possible for a spike initiated in the dendrite to reach the soma before other synapses. This produces a negative delay between the soma and these delayed synapses, even though the spike was dendritically initiated. It is not possible however, for a spike initiated in the soma to show a positive delay. C) Distribution of spike times recorded at all weak pathway synapses. Spike times are shown relative to the onset of the corresponding burst. D) Model prediction comparing plasticity in the weak pathway when it is unpaired (weak only) and paired (weak + strong). The vertical axis (Norm. Weight) is the average weight of all weak pathway synapses at the end of simulation, calculated offline using the voltage-based learning rule [41]. E) Experimental data (same as Fig. 3D) shown again for comparison with panel D here. Both model and experiment show that anodal DCS increases LTP in the weak pathway only when it is paired with strong pathway activation. Bottom row: simulations and methods are identical to the top row, with two exceptions. First, we emulated the application of locally applied somatic TTX by setting voltage-gated sodium conductance to zero in the soma and axon, preventing the initiation of spikes in these compartments. Second, the number of synapses in each pathway was doubled, increasing the likelihood of spike generation, which now occurred in the dendrite. The testable prediction of the model is that in the presence of TTX now DCS will no longer boost LTP.

The results of Fig. 7 highlight the complex interaction between endogenous synaptic input dynamics, synapse location, and DCS-induced polarization. Despite the complexity, Fig. 7 also points to a simple and more general principle: when endogenous plasticity is primarily driven by *somatic* sources of depolarization (e.g.

backpropagating somatic spikes), DCS-induced polarization at the *soma* determines effects on plasticity. This is what we observe with TBS (Fig. 7 bottom two rows). When endogenous plasticity is primarily driven by *dendritic* sources of depolarization (e.g. subthreshold depolarization or dendritic spikes), DCS-induced

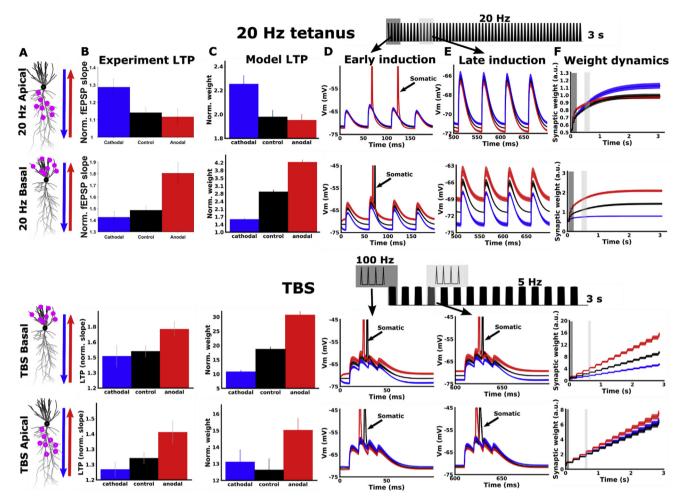


Fig. 7. Model captures interaction between dendritic location and induction protocol. A) Simulated neuron morphology, showing distribution of activated synapses for 20 Hz (top two rows) and TBS (bottom two rows). Arrows indicate the direction of the DC fields for anodal (red) and cathodal (blue) stimulation. B) Experimental LTP results for each condition. The vertical axis is the average of the last ten normalized fEPSP responses. The top two panels are reproduced from data in [14]. The bottom two panels are identical to Fig. 4D, shown again here for comparison. C) Model LTP predictions qualitatively match experimental LTP results (c.f. C; same direction of DCS effect). The vertical axis (Norm. Weight) is the average weight of all activated synapses at the end of simulation, calculated offline using the learning rule of [41]. D) Example simulated voltage traces for individual cells recorded only at activated synapses during the first four input pulses. Traces are averaged over all activated synapses for the example cell. Spikes that back-propagate from the soma are indicated with arrows. E) Same as D, but at a later time point in the simulation (pulses 10–13 for 20 Hz tetanic stimulation; pulses 13–16 for TBS simulations). Note that for 20 Hz stimulation synaptic depolarization is reduced due to short term depression and somatic spiking ceases very early in the simulation. During this subthreshold period, DCS causes a small shift in membrane potential and the resulting plasticity. Since DCS causes opposite subthreshold polarization in apical and basal dendrites, the effect on LTP is also opposite in apical and basal dendrites (C, top two rows). For TBS simulations, recovery from short term depression between bursts allows bursts later in the simulation to produce somatic spikes. Plasticity throughout the simulation is controlled by somatic spikes, and is similar in apical and basal dendrites (C, bottom two rows) F) Dynamics of synaptic weights during the full simulation, averaged over the entire population of activated synapses. For TBS simulations, the weight change is approximately linear in the number of bursts, as each successive burst is equally effective at inducing plasticity. For 20 Hz stimulation, the weight change saturates with the number of pulses, as each successive pulse is weaker due to short term depression. Only the weight at the end of the simulation is used to predict the resulting LTP in experiments (C). Gray boxes in F indicate time periods for early (dark gray) and late induction (light gray) that are plotted in D and E, respectively. A schematic of the input pulse train and relative timing of early (dark gray) and late (light gray) induction pulses are shown at the top. All data are represented as mean ± s.e.m. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

polarization at the *dendrite* determines effects on plasticity. This is what we observe with 20 Hz tetanus (Fig. 7 top two rows) or when we block somatic spiking (Fig. 6 bottom row).

Dose response and distribution of plasticity effects

We are ultimately interested in understanding the effects of weaker electric fields that occur in the human brain during clinical tDCS, which are on the order of 1 V/m [45,46]. The model presented above is able to reproduce several experimental effects of DCS (Figs. 5–7, Supplemental Fig. S3) and canonical synaptic plasticity results (Supplemental Fig. S1) with the same set of parameters (Supplemental Table 1). Because the model includes the actual

morphology of CA1 pyramidal neurons, the electric field magnitude in simulations has a precise mapping to the electric field in experiments. We therefore used the model to make predictions for how weaker electric fields would influence synaptic plasticity.

We first measured the passive membrane polarization throughout the model neuron in response to DCS (Fig. 8A). As observed experimentally [18], we found that the subthreshold membrane polarization is linear in the electric field magnitude, with opposite polarization in the soma and apical dendrites (Fig. 8B). Next we repeated simulations of TBS with DCS at varying electric field magnitudes (+- 1, 2, 5,10,15,20 V/m). For a given electric field magnitude, we quantified the effect of DCS at each synapse and averaged over all active synapses in the apical dendrite

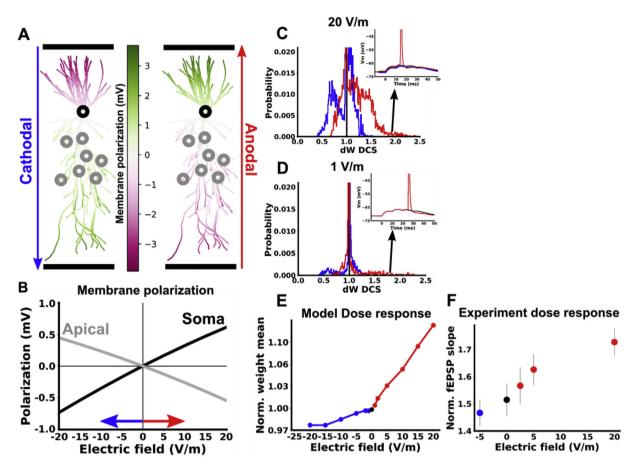


Fig. 8. Dose response in computational model and experiments. A) Membrane polarization of a CA1 pyramidal cell in response to 20 V/m cathodal (left) and anodal (right) electric field. B) Membrane polarization in response to varying electric field magnitude. On the horizontal axis positive values correspond to anodal DCS and negative values correspond to cathodal DCS. The gray curve is averaged over all segments in the apical dendrite, and the black curve is measured at the soma. C,D) Distribution of DCS effects on synaptic weight in response to TBS in apical dendrites. The horizontal axis is the final synaptic weight during a simulation with DCS divided by the final synaptic weight in the same cell under control conditions. dW_{DCS} therefore measures the change in weight caused by DCS for each synapse. Inset shows example voltage traces for synapses in the tail of the distribution. These synapses correspond to cases where the control simulation brought the cell to slightly below threshold, such that DCS was able to cause firing and produce a large change in the weight. E) Mean of the synaptic weight change (dW_{DCS}) due to DCS, averaged over all simulated apical synapses, as a function of DCS electric field. F) Experimental LTP as a function of DCS electric field. All data are represented as mean \pm s.e.m.

(gray circles in Fig. 8A; see Methods). We found that the mean effect of DCS on plasticity is monotonic in the electric field magnitude (Fig. 8E). While each polarity of electric field produces an approximately linear dose response, we observed a greater slope for anodal (positive) electric fields. This asymmetry of anodal and cathodal effects is consistent with our experimental observations (Fig. 1B).

If we consider the distribution of DCS effects over all apical synapses, we find that for weak fields the mean effect of DCS is predominantly driven by the tail of this distribution, where very few synapses have large changes in plasticity (Fig. 8D). For a small number of cells that are close to threshold, a weak field may cause a spike that would have otherwise not happened. This causes a large jump in all synaptic weights for a few highly sensitive cells. While most synapses see very small effects on their weights due to small effects on spike timing and subthreshold polarization, a small number of synapses experience a large effect on their weights due to the initiation of new spikes.

We next experimentally tested the dose response by varying the DCS electric field (-5, 0, 2.5, 5, 20 V/m; cathodal negative; anodal positive). Consistent with the prediction of the model, there was a monotonic relationship between electric field and the magnitude of LTP (Fig. 8F): -5 V/m (1.41 ± 0.049 , N = 16), 0 V/m (1.46 ± 0.060 , N = 20), 2.5 V/m (1.52 ± 0.068 , N = 15), 5 V/m (1.57 ± 0.057 , N = 14),

20 V/m (1.67 \pm 0.051, N = 18), with larger effects for anodal stimulation (see also Fig. 1B). Unlike the model however, the effect of anodal stimulation appears to be saturating at 20 V/m, perhaps reflecting saturation of LTP itself. This discrepancy can be accounted for by considering that synaptic weights in the model can grow without bound, unlike biological synapses [47].

Discussion

Synaptic plasticity is critical for many forms of learning and tDCS has been thought to alter synaptic plasticity [6,48]. How stimulation may interact with ongoing synaptic activity to alter plasticity remains poorly understood. Here we found that weak electrical stimulation with constant direct currents can enhance LTP, while maintaining input specificity and associativity. We propose a model in which DCS boosts endogenous Hebbian synaptic plasticity through modulation of pyramidal neuron membrane potential dynamics. As this model predicts, the effects of DCS also reflect the input specificity and input associativity of the endogenous Hebbian plasticity.

This framework produces a number of testable predictions for clinical experimentation. First, the efficacy of clinical protocols should improve when tDCS is paired with a learning task which induces plasticity, instead of the common practice of pairing tDCS

with "rest". Second, when tDCS is paired with a learning task, we postulate that the effects should be highly specific to the trained task. Finally, the pairing of tDCS with Hebbian plasticity and learning can be thought of as a method for functional targeting, since tDCS should only affect synaptic pathways that are already undergoing plasticity due to the paired task. This may alleviate the prevailing concern that focal stimulation of a desired target in the brain is not possible with transcranial electrical stimulation.

Hebbian plasticity

Hebb originally proposed that coincident pre and postsynaptic firing was required for enhanced synaptic efficacy [49]. Over time the concept of Hebbian plasticity has come to incorporate forms of plasticity that depend on correlated pre and postsynaptic activity variables, regardless of the exact biophysical implementation [50]. While we do not directly measure or manipulate postsynaptic firing here, TBS-induced LTP at CA1 Schaffer collaterals has been shown to be Hebbian in that it depends on pre and postsynaptic activity and exhibits classic Hebbian properties of input specificity and associativity [51]. The synaptic plasticity rule in our model is similarly Hebbian in that plasticity depends on correlated pre and postsynaptic activity in the form of presynaptic spike arrival and postsynaptic membrane voltage [41].

Input specificity

Input specificity is a property of Hebbian plasticity whereby only synaptic inputs that are coactive with the postsynaptic neuron, and presumably relevant for the current task, are strengthened [35]. The computational significance of this specificity has been recognized for some time, as it allows a network of neurons to learn sparse, non-overlapping neural representations [52]. In practice, this is implemented in the brain by molecular machinery which responds to elevated activity specifically at task-relevant synapses [53]. Here we show that DCS enhances LTP in a manner that respects this input specificity. DCS only boosts the strength of synapses that are active and already undergoing endogenous plasticity. Based on this observation, we make two predictions for the optimization of tDCS effects in humans.

First, tDCS effects in humans should similarly exhibit synaptic input specificity, which would be reflected as task specificity in the cognitive domain. Indeed, there is good evidence for task-specific effects of tDCS, despite its lack of spatial focality in the brain [54,55]. This property may be central to the broad application of tDCS. It implies that tDCS can be used flexibly in combination with many different tasks and with limited side effects, despite stimulation reaching large regions of the brain. Second, tDCS effects may be most pronounced when paired concurrently with training that induces plasticity. Again, there is evidence for this in the human literature [33,56]. It may be possible to leverage these properties further by pairing stimulation with forms of learning that are known to rely heavily on Hebbian mechanisms [57–59].

Associativity

Associativity refers to the potentiation of a weak synaptic input when it is paired with strong input at other synapses to the same neuron. In this sense the weak input becomes associated with the strong input. This can serve as a cellular mechanism to bind previously unrelated information as in classical conditioning [60], and to form cell assemblies for associative learning [35–37]. Here we show that anodal DCS can further enhance this associativity, which may manifest as an increased probability of forming associations between stimuli during learning that involves Hebbian plasticity.

We did not explore associativity under cathodal DCS as we saw no effect on LTP for the single pathway experiment, but we cannot rule out an associative effect of cathodal DCS.

Asymmetry

As in our previous work [14,61] and in many tDCS studies [28,62,63], we observe asymmetric results with respect to DCS polarity. Anodal DCS enhanced LTP, while cathodal DCS had no discernible effect with the current sample sizes. This stands in contrast to the symmetric membrane polarization observed with opposing field polarities [18]. Of course, the brain exhibits highly nonlinear responses to changes in membrane voltage, from the level of ion channels to the propagation of activity in a recurrent network. In this sense, it is perhaps not surprising that responses to DCS are nonlinear. However, it remains a crucial topic to understand which sources of nonlinearity are most relevant for DCS, and whether these persist in human tDCS. Below we speculate on some of these potential sources, although we are unable to disambiguate them here, as it is beyond the scope of the current study.

The asymmetry may result from the interaction between DCS effects on different neuronal compartments. For example, during cathodal stimulation, *depolarization* of *apical dendrites* can counteract *hyperpolarization* of *somas* so that there is no reduction in LTP [61](Fig. 8A and B). However, this mechanism cannot explain the asymmetry we observed for LTP in basal dendrites (Fig. 4C and D; bottom row), as the direction of polarization is the same in both basal dendrites and somas (Supplemental Figs. S5A and B). While our model does predict a nonlinear dose response in basal dendrites (Supplemental Fig. S5E), this is more likely due to nonlinear responses of voltage-gated ion channels or the synaptic plasticity molecular machinery.

A nonlinear dose response may also result from the distribution of initial synaptic states in the cell population that we record from. For example, the prior history of the recorded synapses may be such that they are biased towards an increase in strength [64]. Similarly, it could reflect the distribution of cell excitability, such that cells are biased toward an increase in firing. With this in mind, we analyzed the input-output relationship between fEPSP's and population spikes in our baseline recordings. Indeed, we found that our experiments are run near a nonlinearity in this input-output relationship, such that population spiking could be more readily increased than decreased (Supplemental Fig. S4).

Mirroring the asymmetric effect of DC polarity, the effects with respect to phase of AC stimulation was also asymmetric. This suggests that even in the absence of information about the precise timing of synaptic inputs when tACS is applied in humans, a net enhancement of LTP may be expected when tACS is paired with synaptic plasticity induction. Notably, the boost in LTP was also larger here for ACS than DCS, perhaps owing to the frequency response properties of pyramidal neuron membranes showing a peak at theta frequencies [20,65].

Mechanism

Perhaps the most well characterized cellular effect of electrical stimulation is the modulation of somatic membrane potential and firing probability [18,20,61,65–70]. In human tDCS studies, it is the modulation of motor-evoked potentials, which have been linked to long-term plasticity [48,71,72]. Here we propose a model which translates acute changes in firing probability and timing into long term changes in synaptic plasticity. In addition to several other phenomena [7,8,73,74], previous studies have pointed to the effects of DCS on BDNF release [12,13,15,75]. While the precise mechanisms remain unclear, BDNF appears to be released in response to

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postsynaptic depolarization and involved in LTP induction [76–78]. BDNF may therefore be an essential part of the molecular machinery that translates DCS-induced effects on membrane potential dynamics into changes in plasticity.

Electric fields are also known to alter cell motility and immune responses [79,80]. However, these effects unfold over the course over many minutes to hours. During prolonged stimulation, it is likely that various effects on cellular physiology begin to take hold simultaneously, with interactions between them. However, robust effects were generated here with remarkably short stimulation duration (3 s), which depended on stimulation polarity with subsecond timing (100 ms, Fig. 1C). Polarization of neuronal membranes is the only known effect of stimulation that acts on these timescales, making it a likely source of effects here. Inhibitory neurons were not included in our model as the effects of DCS are expected to be small, at least for neurons with symmetric morphology [19]. However, we cannot rule out that DCS polarizes inhibitory neurons on a rapid time scale, which in turn may affect plasticity either directly or through network effects. Prolonged stimulation necessarily includes effects operating on both short timescales (membrane polarization, plasticity induction) and longer timescales (cell motility and immune responses), e.g. [81]. However, shortening the stimulation and pairing it with quicker (sub-minute) bouts of training as we have done here, could be a useful strategy to isolate effects based on Hebbian plasticity induction, which operate on faster timescales.

Our experiments and computer simulations support a model in which DCS affects TBS-induced LTP primarily by somatic polarization and changes in somatic spiking (Figs. 4–6). However, DCS-induced dendritic polarization is also likely to contribute to plasticity, as we suspect for 20 Hz tetanus experiments (Fig. 6) [14]. Our computational model can reconcile these results by considering the voltage dynamics during induction (Fig. 6).

We propose a general principle that emerges from this result: when endogenous plasticity is primarily driven by somatic sources of depolarization (e.g. spikes), DCS-induced polarization at the soma determines effects on plasticity. When endogenous plasticity is primarily driven by dendritic sources of depolarization (e.g. subthreshold depolarization or dendritic spikes), DCS-induced polarization at the dendrite determines effects on plasticity. The relative contribution of somatic and dendritic DCS effects, and therefore the overall effect on plasticity, is not always obvious. The spatial location and temporal pattern of active synapses (including background synaptic input), as well as neuromodulator concentrations and intrinsic excitability can all shift the endogenous voltage dynamics towards somatic or dendritic dominance. Computational models, such as the one presented here, can help in this regard by exploring how DCS interacts with this large parameter space of endogenous synaptic activity. This should be an important next step for future work.

Brain region

While electric current does reach the hippocampus and subcortical structures during stimulation [82], tDCS is thought to primarily act on neocortex. Here we chose hippocampus as a model system for the wealth of studies on hippocampal synaptic plasticity and the much neater organization of input pathways. While not identical, many excitatory plasticity mechanisms are conserved in pyramidal neurons between cortex and hippocampus [83], making our observations here informative for cortex as well. Indeed, the plasticity rule used here in our model has also been used to describe plasticity at neocortical excitatory synapses [26,41,43]. Of course, further work is needed to validate this relationship with respect to DCS effects. It is also worth noting that this work, in

addition to other recent studies [12,15,84], motivates the hippocampus as a target for tDCS.

Dose response

Here we used a 20 V/m electric field in order resolve effects with a reasonable number of animals. Electric fields in the brain during typical tDCS experiments are expected to be 1 V/m or less [45]. While we do not measure effects with this intensity, our computational model predicts a monotonic relationship between the population-mean synaptic plasticity and electric field magnitude (Fig. 8C). For a given DCS polarity, the model predicts a linear relationship between field magnitude and mean plasticity effects (Fig. 8E). To first order this implies population mean effects of ~1% for fields of 1 V/m (we observe ~20% effects for 20 V/m), in line with effect sizes observed for acute effects of DCS [70]. However, experimentally we observe a saturation with increasing stimulation intensity (Fig. 8F). This linear approximation may therefore underestimate effect sizes with weaker fields.

We also note recent efforts to increase stimulation intensity up to 6 mA in humans by distributing current across multiple electrodes [85], which can achieve electric fields of 3 V/m in the brain [82]. Given our estimates here, this would generate effects on synaptic plasticity of at least ~3%, notably affecting a few synapses most strongly (Fig. 8D). Recent in vivo rodent work suggests that a motor learning task leads to potentiation of ~1–2% of synaptic spines in a given volume of cortex [86], which is comparable to what we expect tDCS to achieve. Effect sizes of tDCS on synaptic plasticity in humans are therefore likely to be in a behaviorally relevant range.

Methods

All animal experiments were carried out in accordance with guidelines and protocols approved by the Institutional Animal Care and Use Committee (IACUC) at The City College of New York, CUNY (Protocol 846.3 and 2016–24).

Brain slice preparation

Hippocampal brain slices were prepared from male Wistar rats aged 3-5 weeks old, which were deeply anaesthetized with ketamine (7.4 mg kg^{-1}) and xylazine (0.7 mg kg^{-1}) applied I.P., and sacrificed by cervical dislocation. The brain was quickly removed and immersed in chilled (2-6 °C) dissecting artificial cerebrospinal fluid (aCSF) solution containing (in mM): Choline chloride, 110; KCl, 3.2; NaH₂PO₄, 1.25; MgCl₂, 7; CaCl₂, 0.5; NaHCO₃, 26; D-glucose, 10; sodium ascorbate, 2; sodium pyruvate, 3. Transverse slices (400 µm thick) were cut using a vibrating microtome (Campden Instruments) and transferred to a chamber containing a recovery aCSF at 34 °C: NaCl, 124; KCl, 3.2; NaH₂PO₄, 1.25; MgCl₂, 1.3; CaCl, 2.5; NaHCO₃, 26; D-glucose, 25; sodium ascorbate, 2; sodium pyruvate, 3. After 30 min in the recovery solution, slices were transferred to a holding chamber containing recording aCSF at 30 °C: NaCl, 124; KCl, 3.2; NaH₂PO₄, 1.25; MgCl₂, 1.3; CaCl, 2.5; NaHCO₃, 26; D-glucose, 25; for at least 30 min. Finally, slices were transferred to a fluid-gas interface chamber (Harvard Apparatus) perfused with warmed recording aCSF (30.0 ± 0.1 °C) at 2.0 ml min⁻¹. Slices were allowed to acclimate to the recording chamber for at least 30 min before recording started. The humidified atmosphere over the slices was saturated with a mixture of 95% O2-5% CO2. All aCSF solutions were bubbled with a mixture of 95% O_2 –5% CO_2 . Recordings started approximately 2 h after the animal was sacrificed.

fEPSP recordings

Field excitatory postsynaptic potentials (fEPSP's) were evoked using a platinum-iridium bipolar stimulating electrode placed in either stratum radiatum (apical experiments) or stratum oriens (basal experiments) of CA1 within 200 μm of the somatic layer. Recording electrodes were made from glass micropipettes pulled by a Sutter Instruments P-97 and filled with recording aCSF (resistance 1–8 M Ω). A "dendritic" recording electrode was placed in stratum radiatum (apical) or stratum oriens (basal) approximately 400 µm from the stimulating electrode in CA1 to record fEPSP's. The stimulating electrode and dendritic recording electrode were placed at approximately the same distance from the CA1 somatic layer. For all experiments, a second "somatic" recording electrode was placed in the CA1 somatic layer to record population spikes. For two-pathway experiments (Figs. 2 and 3), a second stimulating electrode was placed on the opposite side of the dendritic recording electrode.

fEPSP's were quantified by the average initial slope, taken during the first 0.5 ms after the onset of the fEPSP. The bipolar stimulus intensity was set to evoke fEPSP's with 30-40% of the maximum slope, which was determined at the onset of recording. Baseline fEPSP's were recorded once a minute for at least 20 min before any plasticity induction was applied and only if a stable baseline was observed. For two pathway experiments, stimulation of each pathway was interleaved with an offset of 30 s. After plasticity induction, fEPSP's were again recorded once per minute for 60 min. To measure synaptic plasticity, all fEPSP slopes were normalized to the mean of the 20 fEPSP's immediately preceding induction. The amount of LTP in each slice is quantified as the mean of the last 10 min of normalized responses (51–60 min after induction). Group data are reported as mean and standard error of the mean across slices. Number of slices is indicated with variable N wherever statistics are reported.

DCS

Uniform extracellular electric fields (±20 V/m) were generated by passing constant current (D/A driven analog follower; A-M Systems, WA, USA) between two large Ag-AgCl wires (1 mm diameter, 12 mm length) positioned in the bath across the slice starting 0.5 s before the onset of TBS and ending 0.5 S after the end of TBS (4s total). Slices were oriented such that the somatodendritic axis of CA1 pyramidal neurons was parallel to the electric field between the DCS wires (Fig. 1A). We name each polarity of DCS based on the orientation of the field relative to CA1 pyramidal neurons, and how pyramidal neurons are expected to be polarized. Here, anodal DCS depolarizes CA1 pyramidal neuron somas as it is expected to do in cortical pyramidal neurons under an anode in tDCS. Cathodal stimulation refers to the opposite polarity. Before each recording, DCS current intensity was calibrated to produce a 20 V/m electric field across each slice (typically 100-200 μA) by adjusting the current so that two recording electrodes separated by 0.8 mm in the slice measured a voltage difference of 16 mV (16 mV/ 0.8 mm = 20 V/m).

Quantifying population somatic activity

To quantify the amount of somatic activity in response to synaptic input we used the following method. Raw voltage data recorded in the somatic layer was filtered with a 300 Hz highpass ARMA filter. The filter was designed using the butterworth algorithm via the signal.iirdesign function in the scipy package (design parameters: fs = 10000; fs = 1

defined somatic activity for each evoked response, *s*, as the integral of the high frequency envelope:

$$s_{ij} = \int_{t_{ij}+2ms}^{t_{ij}+8ms} dt \; env(t)$$

$$env(t) = \left| H\left(x_{soma}^{high}(t)\right) \right|$$

where x_{soma}^{high} is the highpassed extracellular voltage recorded in the somatic layer, and H() is the hilbert transform calculated in python using signal.hilbert from the scipy package. t_{ij} is the onset time for the i^{th} input of the j^{th} burst, where $i \in \{1,2,3,4\}$ and $j \in \{1,2,...15\}$. The somatic activity was calculated as the integral of this high frequency envelope in the time window 2–8 ms after t_{ij} , chosen to avoid including the bipolar stimulus artifact. Somatic activity was then normalized to the mean of baseline values (20 responses prior to induction). The same method was used to calculate somatic activity in the population of model neurons (Fig. 5C), except the recorded extracellular voltage in the somatic layer was replaced with the intracellular somatic voltage averaged over all simulated model neurons.

Quantifying population spike timing

We expected that DCS would cause a shift in the average spike timing in the population during TBS (Supplemental Fig. S2E). To create a measure of the mean spike timing, we performed a center of mass calculation on the somatic activity envelope

$$\overline{t_j} = \sum_i \frac{1}{s_{ij}} \int_{t_{ij}+2ms}^{t_{ij}+8ms} dt \ t \cdot env(t)$$

where $\overline{t_j}$ is the population spike timing of the j^{th} burst, env is the envelope of the highpassed extracellular voltage (see above), t_{ij} is the onset time for the i^{th} input of the j^{th} burst, where $i \in \{1,2,3,4\}$ and $j \in \{1,2,...15\}$, and s_{ij} is the somatic activity as in the previous section. Again, we restrict the integrals to between 2 and 8 ms after each input pulse to avoid contributions of the bipolar stimulus artifact. $\overline{t_j}$ can be thought of as the temporal center of mass of the somatic activity during a burst. If more neurons in the population fire earlier during the j^{th} burst, then $\overline{t_i}$ should decrease.

Quantifying dendritic integration

To quantify the amount of dendritic integration in response to synaptic input we used the following method. Raw voltage data recorded in the dendrites (either stratum radiatum or stratum oriens) was filtered with a $5-50\,\mathrm{Hz}$ bandpass ARMA filter. The filter was designed using the butterworth algorithm via the signal-iirdesign function from the scipy package (parameters: fs = 10000; nyquist = fs/2; wp = [5/nyquist, 50/nyquist]; ws = [0.1/nyquist, 100/nyquist]; gpass = 1; gstop = 20; ftype = 'butter'). We then defined dendritic integration, d, for each burst during TBS as the integral of the band-passed signal:

$$d_{j} = \int_{t_{j}+2ms}^{t_{j}+100ms} dt \left| x_{dend}^{low}(t) \right|$$

where x_{dend}^{low} is the band-passed extracellular voltage recorded in the dendrite. For each evoked burst, j, the dendritic integration was

calculated as the integral of this low frequency signal in the time window $2-100\,\mathrm{ms}$ after the onset of the burst, t_j . Dendritic integration was then normalized to the mean of baseline values calculated for each fEPSP (20 responses prior to induction). The same method was used to calculate dendritic integration in the population of model neurons (Fig. 5D), except the recorded dendritic extracellular voltage was replaced with the intracellular voltage averaged over all recorded dendritic segments in the simulated population of cells.

Neuron model

Individual pyramidal cells were modeled in Python using the NEURON simulation package [87]. To construct the model neuron, we reproduced the detailed biophysical neuron model of Migliore et al. [38], and then added parameter changes based on more recent studies. Unless otherwise specified, parameters are the same as in [38].

An L-type calcium channel was added throughout the cell as in [88]. Sodium conductance in the axon was increased to replicate spike initiation in the axon initial segment [89]. Synapses were set to have both AMPA and NMDA receptors, which were modeled as the difference of two exponentials. NMDAR conductance was modified by a voltage dependent mechanism as done previously [88,90,91]. See supplemental tables for the full set of NEURON model parameters.

Synaptic conductances were modified by presynaptic short-term plasticity model as in [92]. Specifically, AMPAR and NMDAR conductances were multiplied by a factor A, which captures short-term facilitation and depression dynamics at presynaptic terminals. A is the product of a facilitation variable F, and 3 depression variables D_1 , D_2 , D_3

$$A = A_0 F D_1 D_2 D_3$$

where A_0 is a constant parameter, which we set to 1 at the start of simulations. At the time of each presynaptic spike, D is multiplied by a factor d such that

$$D \rightarrow Dd$$

while a factor f is added to F such that

$$F \rightarrow F + f$$

Both F and D decay exponentially back towards 1 between spikes according to

$$\tau_D \frac{dD(t)}{dt} = 1 - D(t)$$

$$\tau_F \frac{dF(t)}{dt} = 1 - F(t)$$

Each depression variable D_1 , D_2 , D_3 follows the same dynamics, but with different parameters τ_D and d. The parameters of this short-term plasticity model were fit to the dynamics of fEPSP slopes during the various plasticity induction protocols in [14](i.e. 0.5, 1, 5, 20 Hz trains) and during TBS in this study. The fit was constructed to minimize the squared error between values of A and the normalized fEPSP during induction using the lsqcurvefit function in matlab. See Supplemental Table 2 for the resulting parameters.

The response of an individual pyramidal neuron to the bipolar stimulus in our brain slice experiments was modeled by randomly selecting a group of dendritic segments. AMPAR and NMDAR conductances were then activated simultaneously in the selected

segments. In our experiments we expect that the bipolar stimulus will elicit this synaptic input in a population of pyramidal cells, with the number and location of synapses that are activated varying between cells.

For simplicity, we assume that an integer number of synapses ranging from 5 to 16 can be activated on each cell. This range was selected empirically so that the mean number of synapses produced somatic responses that were close to firing threshold during simulation of TBS. For each integer number of activated synapses, the synapses are randomly distributed on the dendrites, and this was repeated 20 times independently to create a population of 12*20=240 cells. For each cell, synapse locations were drawn randomly and with replacement from a uniform distribution over all dendritic segments that are allowed by the given experiment (e.g. basal dendrites or apical dendrites within 300 μ m from the soma). By sampling with replacement, we allow multiple synapses to be activated on the same dendritic segment, mimicking the random activation of clustered synaptic inputs.

The electric field during DCS was modeled as uniform extracellular voltage gradient. The extracellular voltage at each point in space is then conveyed to each segment of the neuron by NEURON's extracellular mechanism, as has been done previously [23]. Since we are interested primarily in the effect of the extracellular field, for each simulation that applies an electric field there is a corresponding control simulation in which the NEURON model is identical except for the extracellular applied voltage. The effect of the applied field can therefore be compared to a precise counterfactual, where all other aspects of the model are identical.

Voltage-based long-term plasticity rule

We are interested in how synaptic input and postsynaptic voltage dynamics during induction leads to long-term synaptic plasticity (and how DCS can modulate this plasticity). To simulate long-term synaptic plasticity in the model, we use the voltage-dependent plasticity rule of Clopath et al. [41]. As done previously, we assume that actual changes in long-term synaptic strength are delayed relative to the induction period and do not contribute to the dynamics during induction. The plasticity rule is therefore used as a method to compute the final weight change expected at the end of induction and this calculation was done "offline", after simulating the induction protocol. The synaptic weight change is calculated with the following rule (see [41] for further details), which requires information that is derived solely from presynaptic input arrival times and postsynaptic membrane potential measured locally at the synapse:

$$\begin{aligned} \frac{d}{dt}w_i(t) &= -A_{LTD}X_i(t)[\overline{u_-}(t) - \theta_-]_+ + A_{LTP}\overline{x_i}(t)[u(t) - \theta_+]_+[\overline{u_+}(t) \\ &- \theta_-]_+ \end{aligned}$$

where w_i is the weight of the i^{th} synapse, A_{LTD} is a parameter that controls the rate of long-term depression (LTD), A_{LTP} is a parameter that controls the rate of LTP, X_i is the presynaptic spike train, $\overline{x_i}$ is a lowpass filtered version of the presynaptic spike train, u is the postsynaptic membrane potential measured locally at the synapse, $\overline{u_-}$ and $\overline{u_+}$ are lowpass filtered versions of the postsynaptic potential, θ_- is an LTD threshold parameter, θ_+ is an LTP threshold parameter, and $[\cdot]_+$ indicates positive rectification. The dynamics of $\overline{u_-}$, $\overline{u_+}$, and $\overline{x_i}$ are given by

$$\tau_{-}\frac{d}{dt}\overline{u_{-}}(t) = -\overline{u_{-}}(t) + u(t)$$

$$\tau_{+}\frac{d}{dt}\overline{u_{+}}(t) = -\overline{u_{+}}(t) + u(t)$$

$$\tau_{X}\frac{d}{dt}\overline{X_{i}}(t) = -\overline{X_{i}}(t) + X_{i}(t)$$

where τ_- , τ_+ , and τ_x are time constants. We note that in the original learning rule of Clopath et al. [41], A_{LTD} is a function of time and postsynaptic voltage, i.e. $A_{LTD}(t,u)$, which implements homeostatic plasticity. Because we compute synaptic weight changes offline, a homeostatic mechanism is not needed to stabilize the voltage dynamics. To reduce the number of parameters of our model, we therefore treat A_{LTD} as a constant. We apply a lower bound to all synapses such that w_i is set to zero if w_i crosses zero from above.

Parameters for the plasticity model were manually selected so as to replicate classic spike-timing dependent plasticity experiments (Supplemental Fig. S1) and to qualitatively reproduce the effects of DCS on LTP. We are mainly concerned with relative changes in LTP due to DCS (or spike timing/frequency in the case of replicating STDP experiments) and so do not adjust parameters to quantitatively reproduce the amount of LTP in each experiment. Under these constraints we were able to use the same set of parameters for each simulation (Supplemental Table 1). Numerical integration using the forward euler method (0.025 ms time step) was used to solve for $w_i(t)$.

Additional simulation details for Fig. 6

To emulate the two-pathway experiments of Fig. 3, a population of cells was generated as described above, but now two groups of synapses were selected to be activated on each cell, a strong and a weak group. Note that because synapses were selected randomly with replacement, a given synapse was allowed to be part of both groups, although this was rare. As in our experiments, three sets of simulations were run: activation of only the weak pathway at 5 Hz, activation of only the strong pathway with TBS, or activation of both pathways simultaneously.

We hypothesized that pairing the two pathways boosted LTP by facilitating spikes that backpropagate from the soma to synapses in both pathways. To test this hypothesis in our model we wanted to measure spikes that occurred at each synapse, and importantly whether a given spike was initiated in the soma. Of course, if a spike is initiated in the soma, it should occur before a spike is observed in the dendrite. To evaluate this time difference, we first detected the onset of spikes in each segment of the model neuron by measuring the time at which the voltage crossed a threshold of -30 mV from below.

For each segment, a binary vector of spike times was therefore created, with each entry corresponding to a time step in the simulation (1 = spike detected, 0 = no spike detected). The cross correlation was computed between this spike vector and the corresponding vector measured at the soma. This yields binary vector for each segment, where each entry corresponds to a possible time delay between that segment and the soma. A value of 1 in this vector indicates that the corresponding delay was observed. By averaging this cross correlation over all activated synaptic locations, we get a probability density over different spike delays between the soma and dendrite. In general, a spike can propagate throughout the entire neuron within ~2 ms. We therefore assume that temporal correlations occurring within this ± 2 ms window correspond to delays that are due to the propagation of a single spike, while correlations that are outside of this ± 2 ms window correspond to delays between different spikes. We have set up the analysis so negative time delays correspond to spikes that appeared in the soma first. Spikes that initiate in the soma and propagate to the dendrite should therefore add density between -2 and 0 ms (Fig. 6B).

The metric based on spike cross-correlation only captures spike events that occur in both the soma and dendrite. However, spikes can also initiate in the dendrite, but may not propagate completely to the soma. These local spikes would also make a large contribution to synaptic plasticity at a subset of local synapses, but do not contribute to the cross correlation metric. We therefore also considered the overall number of dendritic spikes (global and local) as a function of time during each theta burst at which they occurred. We divided the simulation into individual theta bursts, and within each burst, the simulation was divided into 1 ms time bins. Spikes were counted in these time bins across all synapses. By summing across all synapses, we get the total number of dendritic spikes that occur as a function of the time since burst onset (Fig. 6C).

Simulation details for Fig. 8

Membrane polarization (Fig. 8A and B) was calculated by simulating a single cell without synaptic input for 100 ms with varying applied electric field. Membrane polarization due to DCS was calculated in each compartment as the voltage at the end of the simulation minus the corresponding voltage in the control simulation without DCS.

For each simulation and each activated synapse k, we calculate the effect of DCS on plasticity

$$\Delta W_{DCS}^{k} = \frac{W_{DCS}^{k}(T)}{W_{control}^{k}(T)}$$

where T is the duration of the simulation, $W^k_{DCS}(T)$ is the final weight of the k^{th} synapse at the end of the simulation with DCS, $W^k_{control}(T)$ is the weight at the end of the corresponding control simulation where no DCS was applied. Note that all DCS simulations have a control simulation in which all other details are identical. Therefore, any deviation of ΔW^k_{DCS} from 1 represents the effect of DCS on the k^{th} synaptic weight.

For a given DCS waveform (polarity and magnitude), we are interested in the distribution of ΔW_{DCS}^k over all k synapses in the population. Fig. 8E displays the mean of this distribution as a function of DCS intensity.

Author contributions

G.K. and A.R. designed and executed experiments, M.S. executed experiments, G.K. designed and ran simulations, G.K. analyzed data, G.K., A.R., M.B., L.P. interpreted results and wrote the manuscript.

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Declaration of competing interest

LCP and MB have shares in Soterix Medical Inc and are listed as inventors in patents of the City University of New York related to high-definition tDCS.

Appendix A. Supplementary data

Supplementary data to this article can be found online at https://doi.org/10.1016/j.brs.2019.10.014.

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